



# DYNAMICS OF A 2D GEL AROUND ITS PERCOLATION TRANSITION

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# Aknowledgments

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# Summary

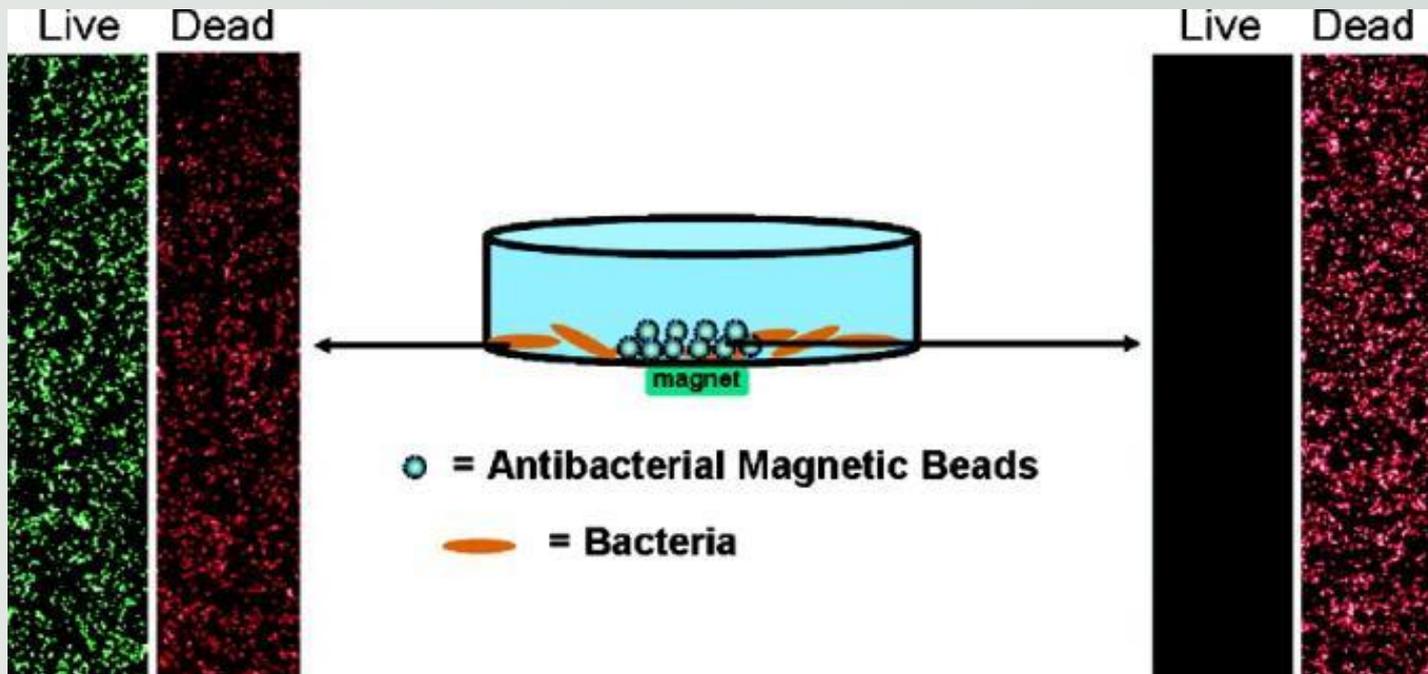
- MOTIVATIONS
- SAMPLE PREPARATION
- **STRUCTURE: IN-SITU MICROSCOPY**
- **MECHANICAL RESPONSE: INTERFACIAL RHEOLOGY**
- **INTERNAL DYNAMICS: XPCS**

# Nanoparticles at interfaces: applications

## Antibacterial coatings

Ag nanocrystal coating of planar surfaces and colloidal particles – antibacterial activity by release of Ag ions.

[ Lee, Cohen, Rubner, *Langmuir* 21, 9651-9 (2011)]



# Nanoparticles at interfaces: applications

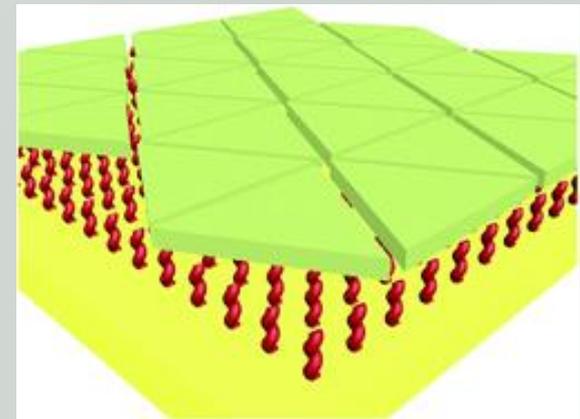
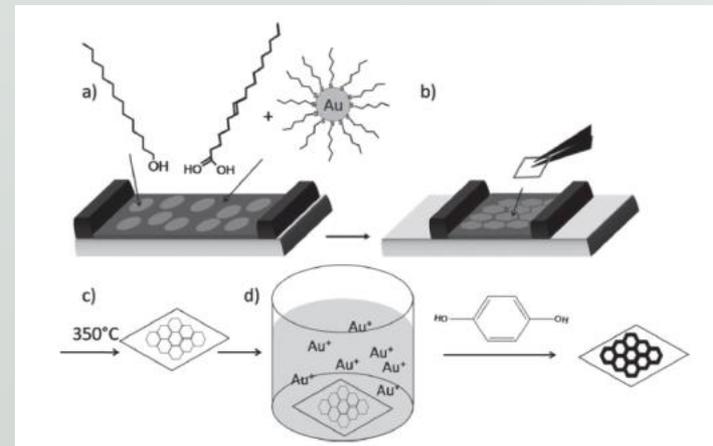
## Electronics and optoelectronics

Thin, large and inexpensive Transparent Conducting Electrodes by Langmuir-Schaeffer Deposition of monolayers of gold nanoparticles

[A. Morag et al. *Advanced Materials* 23, 4327–4331 (2011)]

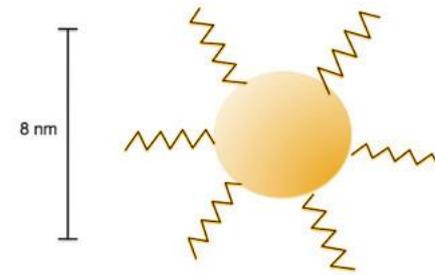
Langmuir Schaeffer deposition of SERS active planar assemblies of gold nanoparticles

[Lee et al. *Nanoscale*, 5, 6404-6412 (2013)]

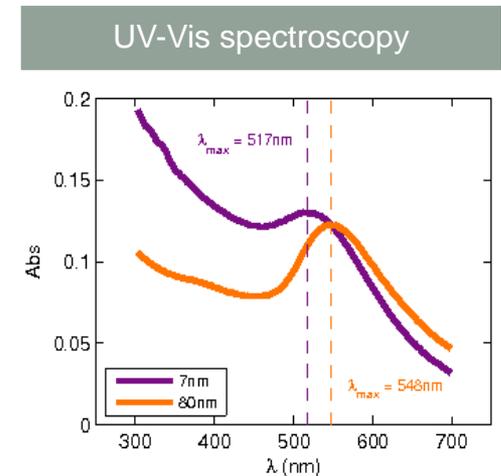
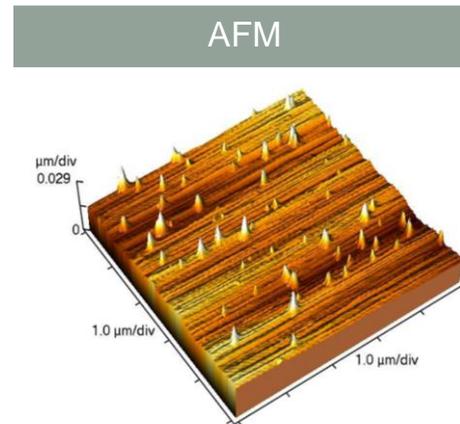
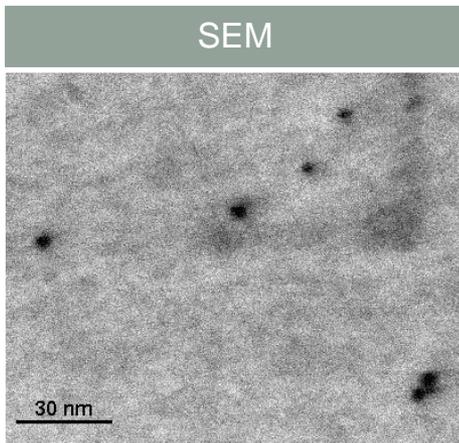


# Our system

Gold nanoparticles (GNPs)  
stabilized by dodecanethiol coating

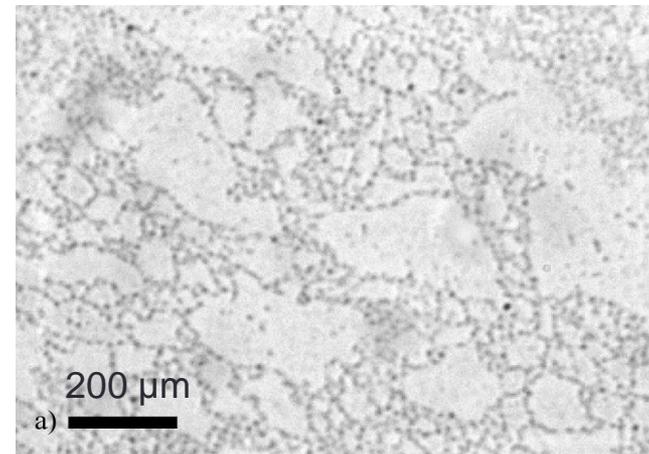
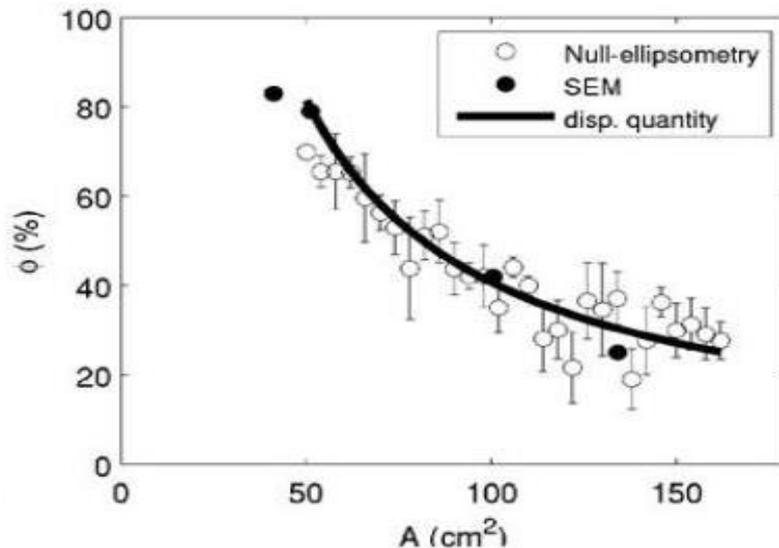


- Synthesized following [Chen, Kimura, *Langmuir* 15, 1075 (1999)]
- 8nm diameter
- Characterized by means of:



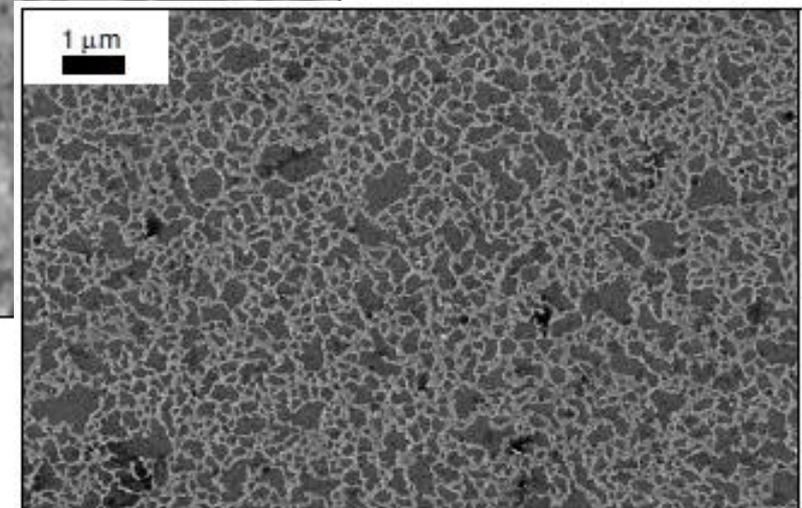
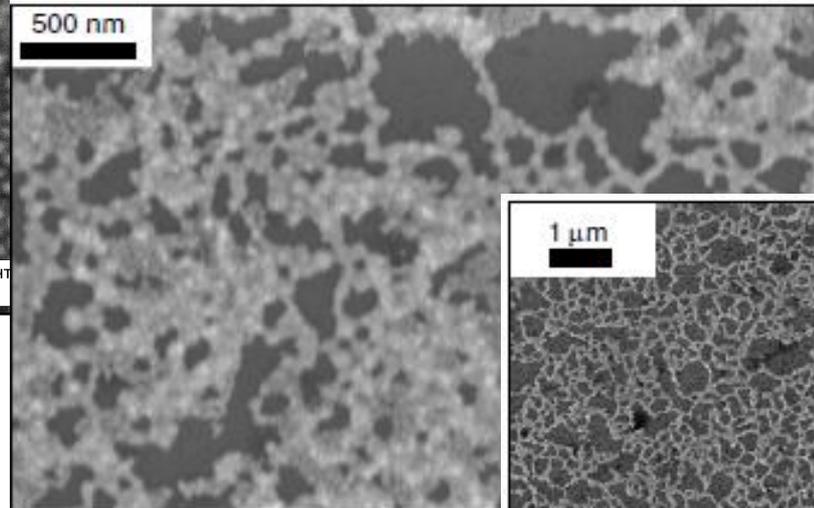
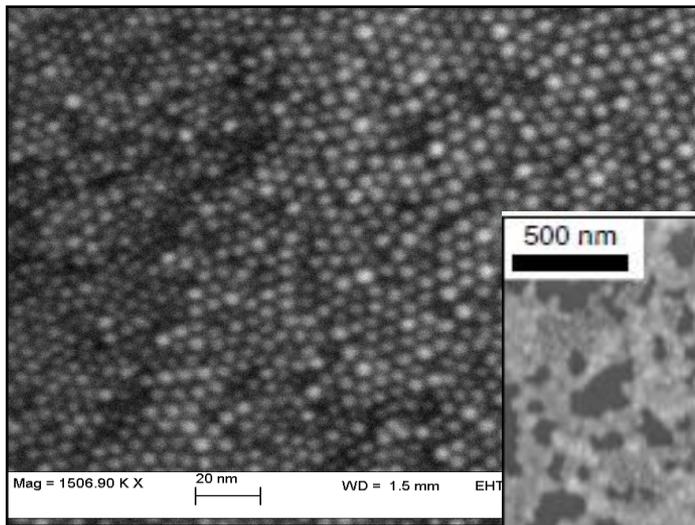
# Langmuir films: preparation

- Film prepared by SLOW spreading of hexane solution at the air/water interface (Langmuir trough)
- The concentration is slowly increased by reducing the available area (trough barrier motion)
- Surface concentration  $\Phi$  of GNPs measured by
  - spreaded amount
  - SEM imaging
  - In-situ microscopy
  - Imaging ellipsometry



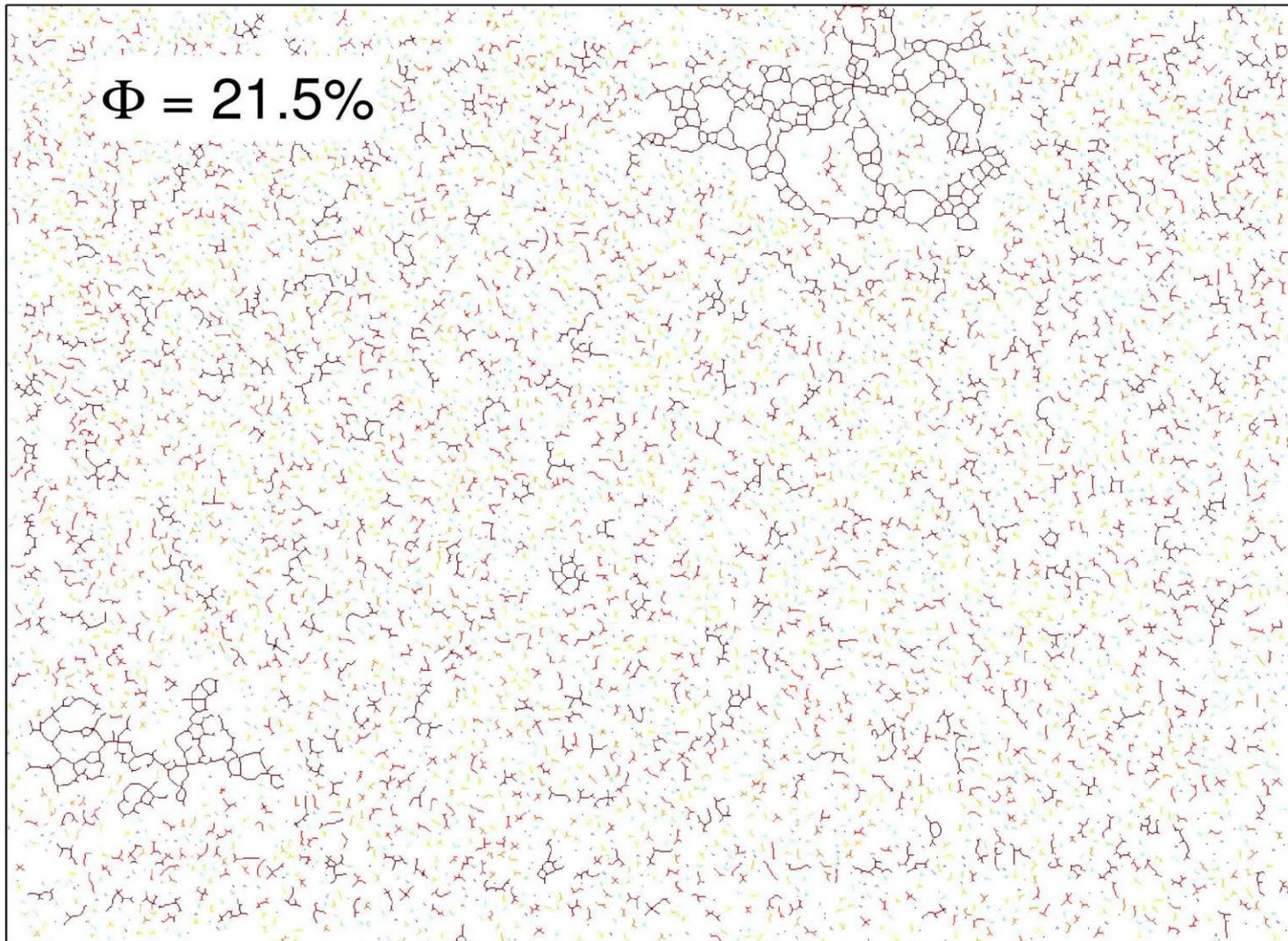
# Langmuir films: SEM Characterization

Transferred film (Langmuir Schaeffer) on silicon substrate

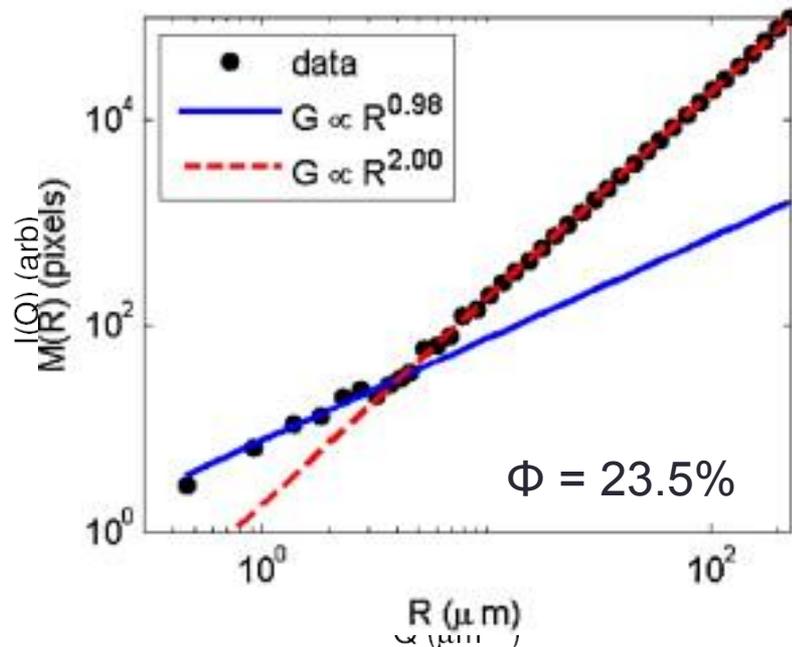
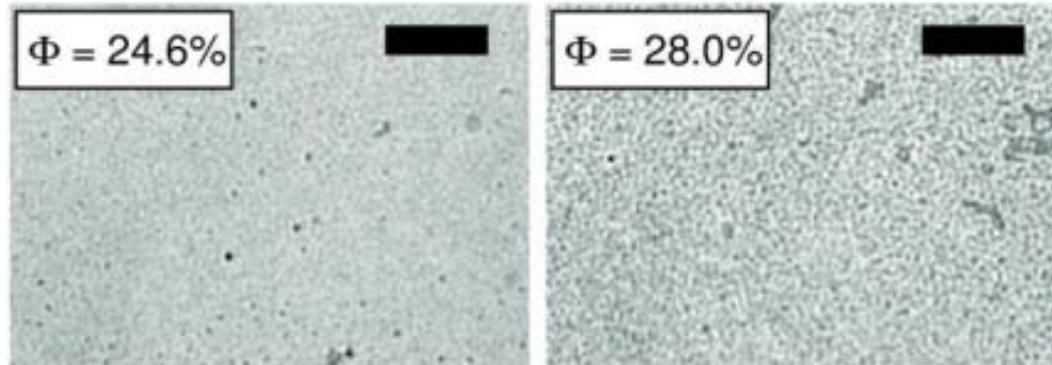


Zeiss SUPRA 40

# Gel formation: in-situ microscopy

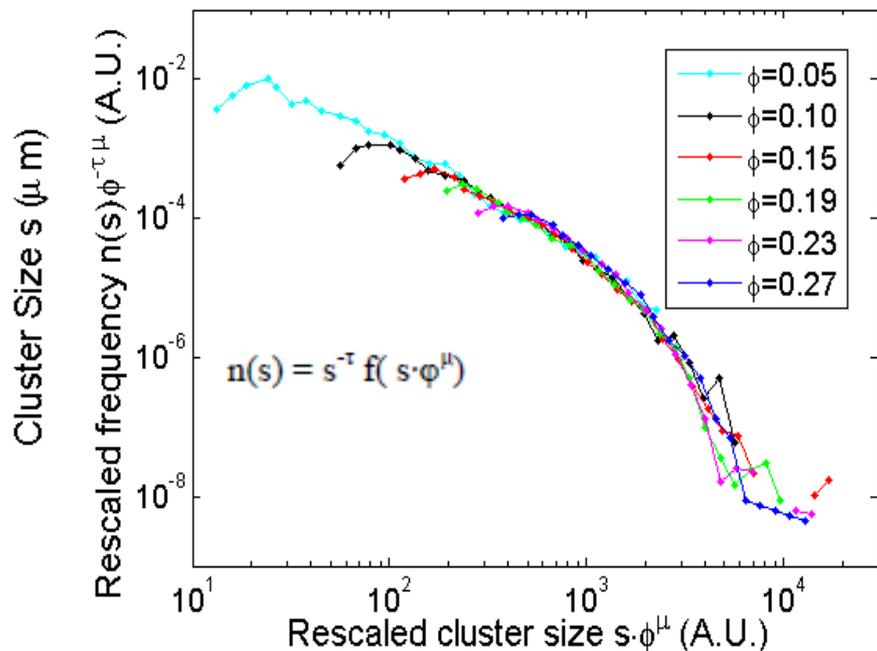
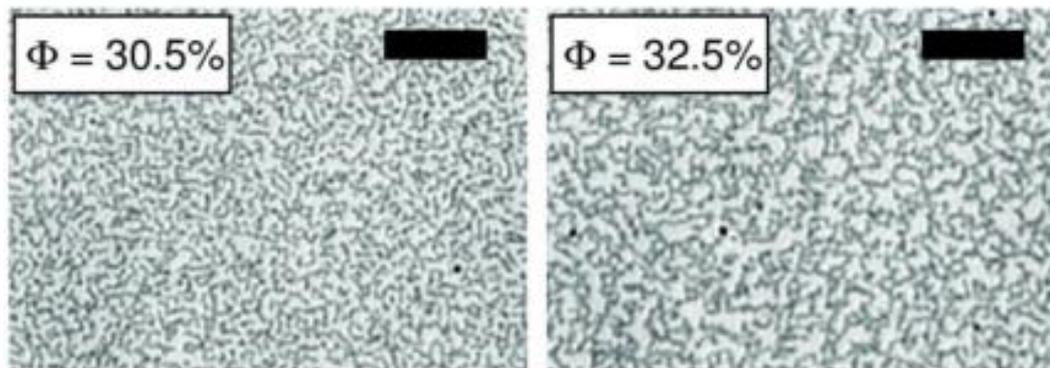


# Gel formation – 1<sup>o</sup> step: flocculation



- crossover (around 4–5  $\mu\text{m}$ ) between a linear and a quadratic growth of the mass
- Small and quasi 1D objects organize in a large scale homogeneous 2D film

# Gel formation – 2<sup>o</sup> step: percolation

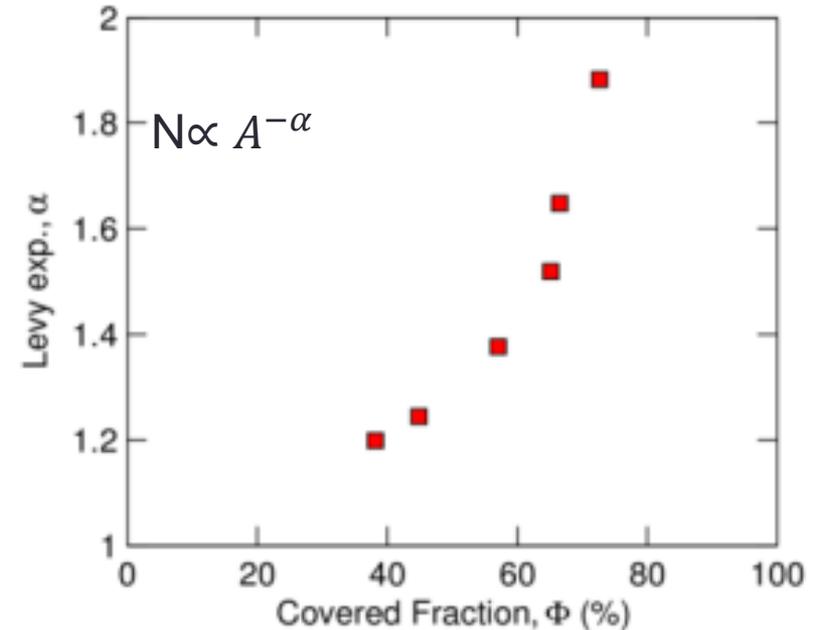
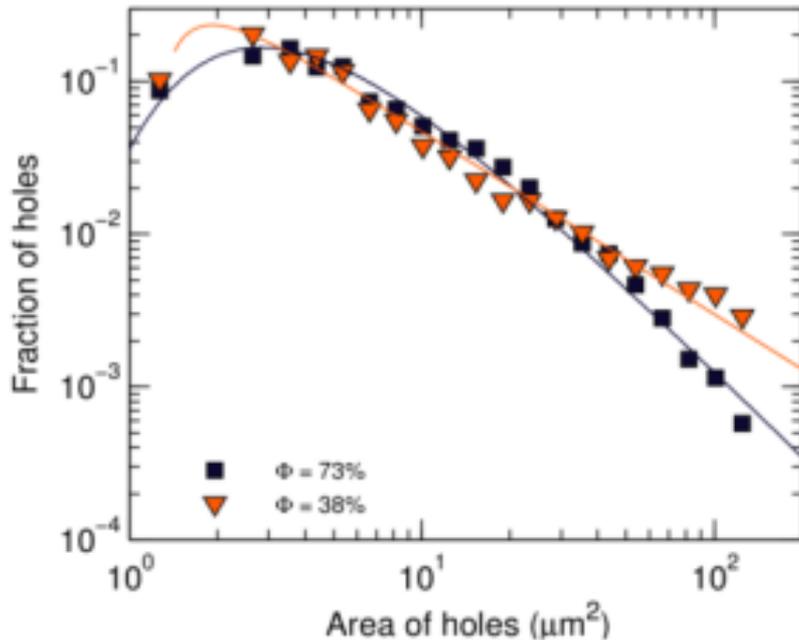
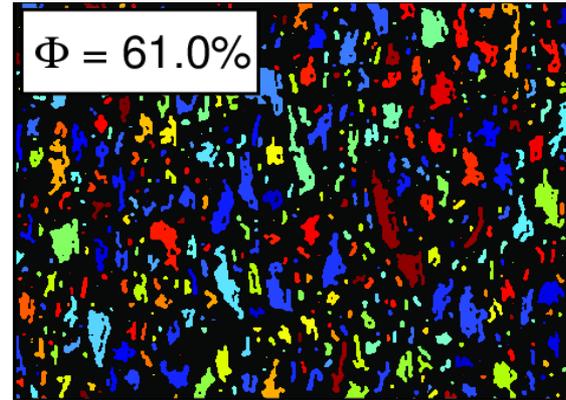
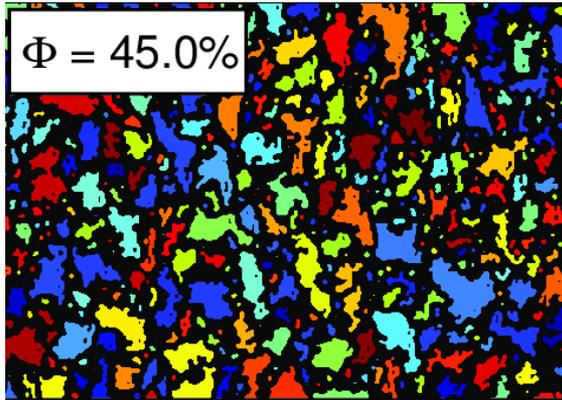


The size distribution scale in a fashion similar to percolation theory

EXPERIMENTS:  $\mu = 1.8 \pm 0.2$      $\tau = 1.6 \pm 0.2$   
 THEORY:             $\mu_c = 91/36 \approx 2.5$      $\tau_c = 187/91 \approx 2.05$

The exponents are **not compatible** with the analytical ones expected for 2D percolation

# Gel formation – 3<sup>o</sup> step: compaction



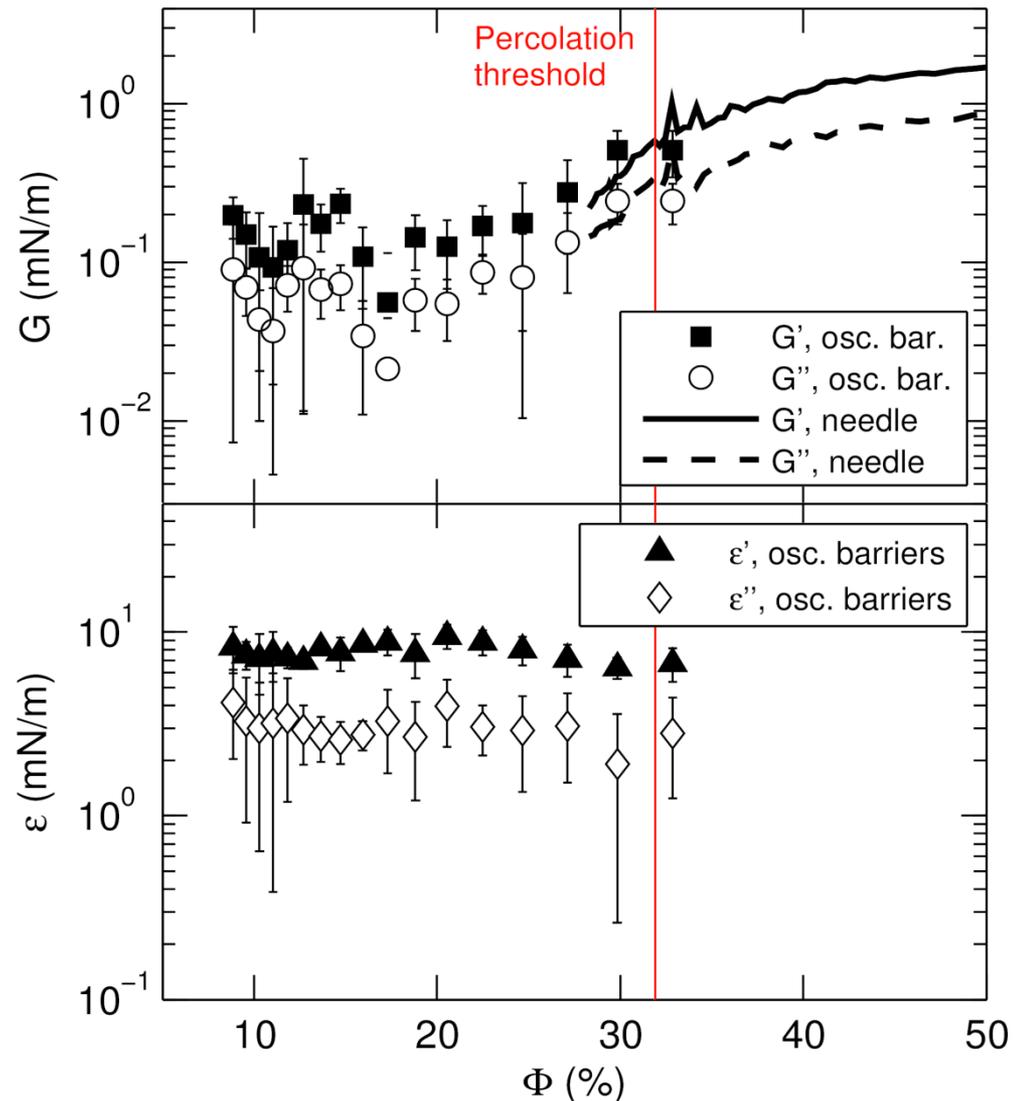
# The gel shows elastic response

Shear  $G$  and compression  $\varepsilon$  moduli measured by means of 2 techniques:

- Oscillating needle
- Oscillating barriers

We find:

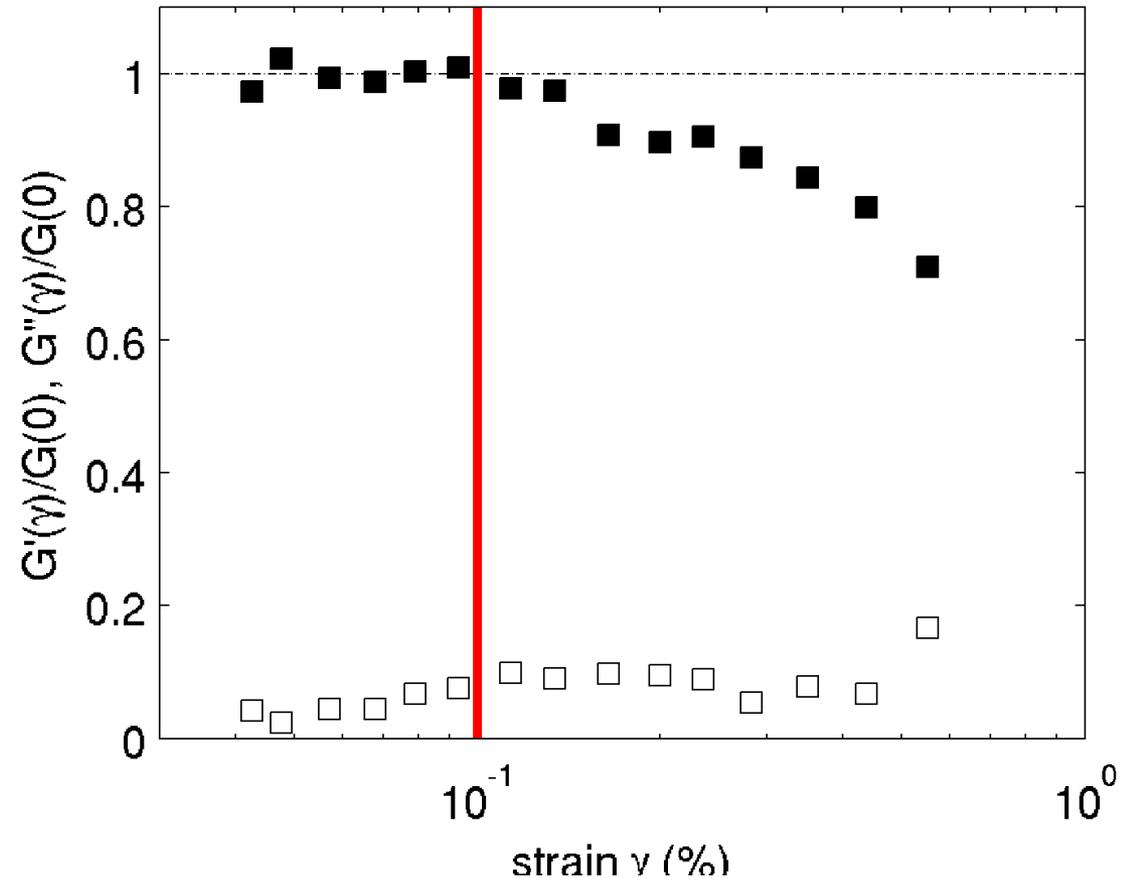
- An mainly elastic response ( $G' > G''$ ,  $\varepsilon' > \varepsilon''$ )
- The compression moduli are constant in the probed  $\Phi$  range
- $G$  increases approaching the percolation transition



# Shear softening

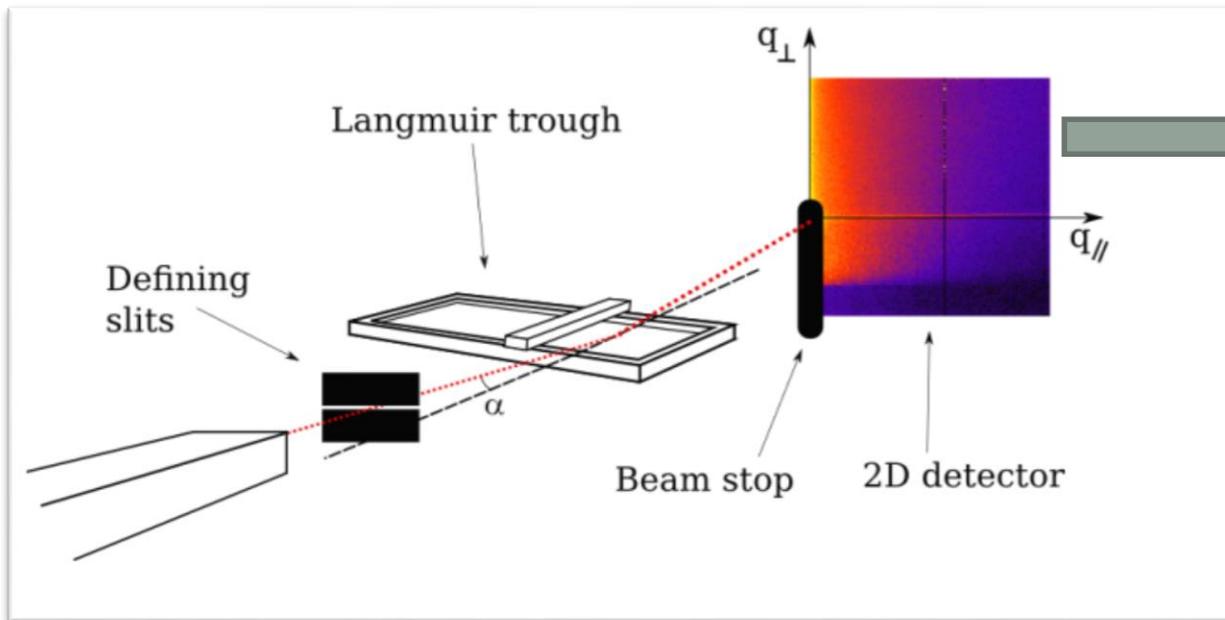
The network structure is very brittle: we have shear softening even at  $\gamma < 0.1\%$

The breakage of weak links modifies the system's structure and reduces the mechanical modulus



# XPCS at the air/water interface

- 8keV coherent beam from the ID10 beamline of ESRF
- SAXS @ Grazing Incidence on water surface ( $\alpha=0.1^\circ$ )
- 2D detector (Medipix, 256x256 pixels) covering
  - q range 0.01 - 0.07 nm<sup>-1</sup>
  - time window 0.01s - 1000s

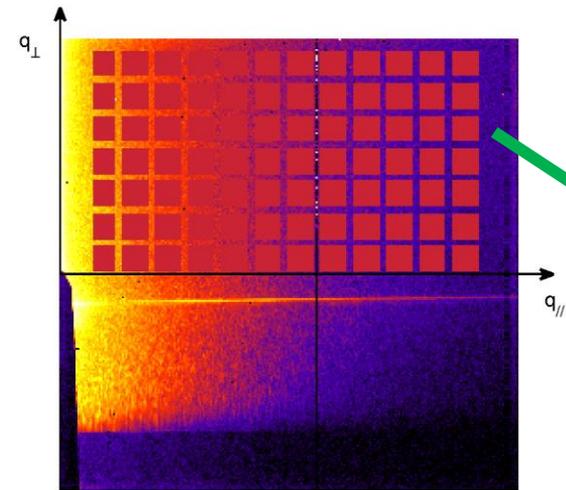


Pixels are grouped in squares identified by  $(q_{\parallel}, q_{\perp})$  coordinates

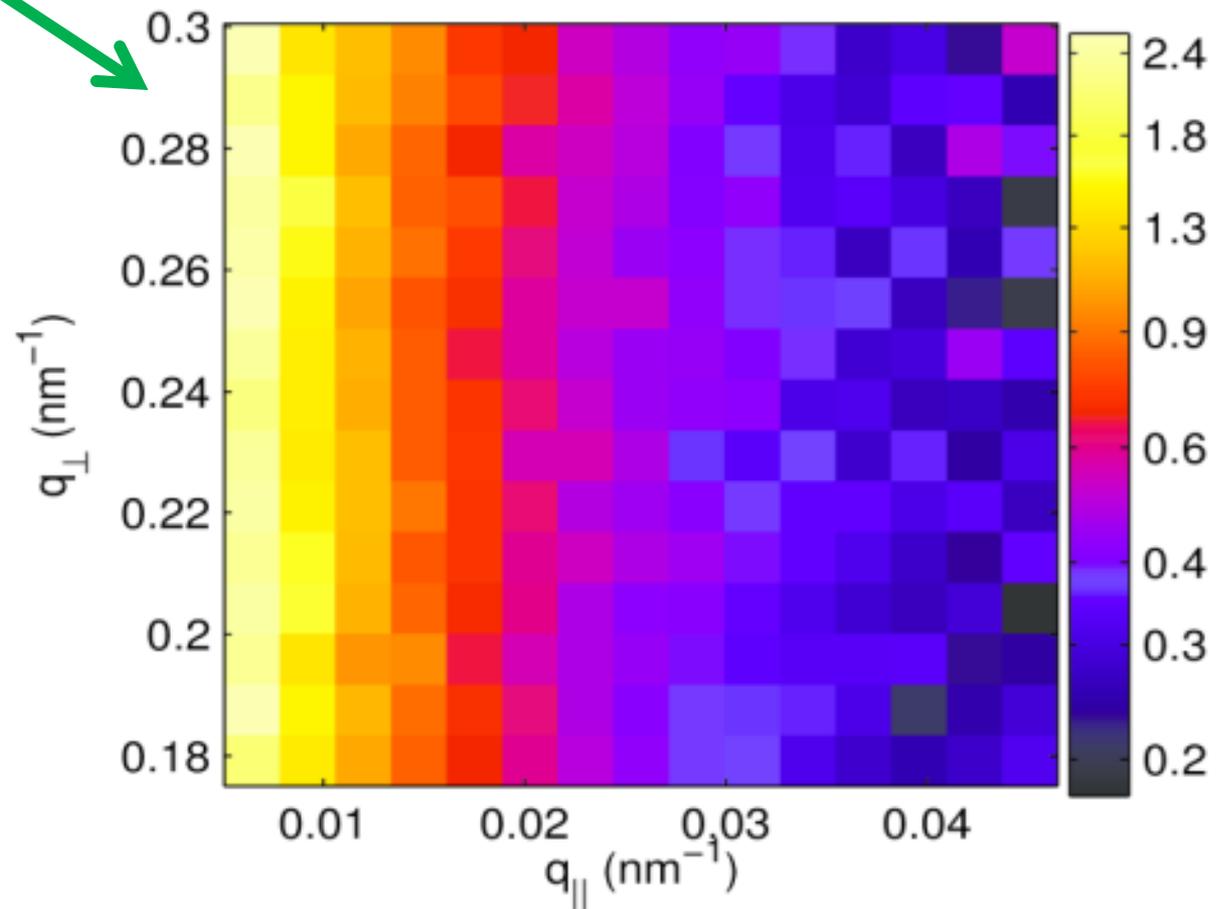
$$g^{(2)}(q, t) = \frac{\langle I(q, t_1)I(q, t_1 + t) \rangle}{\langle I(q, t_1) \rangle^2}$$

$$g^{(2)}(t) = A + \beta e^{-2(t/\tau)^{\gamma}}$$

# Dynamics confined in 2D



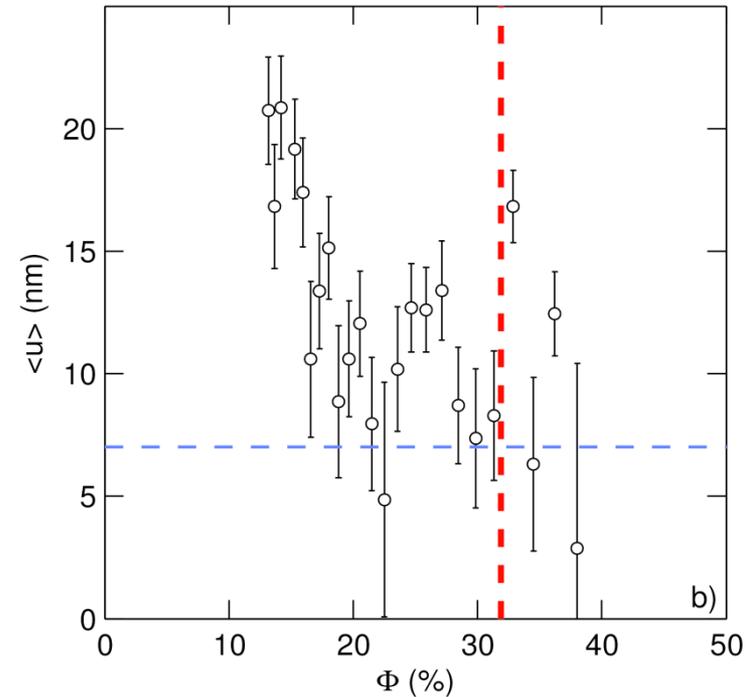
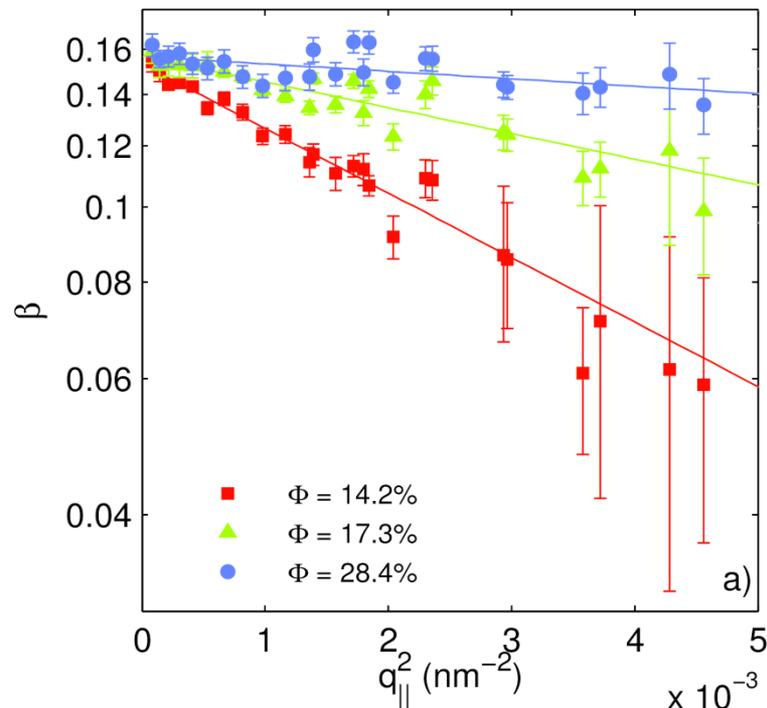
Relaxation time measured at  $\Phi = 15\%$



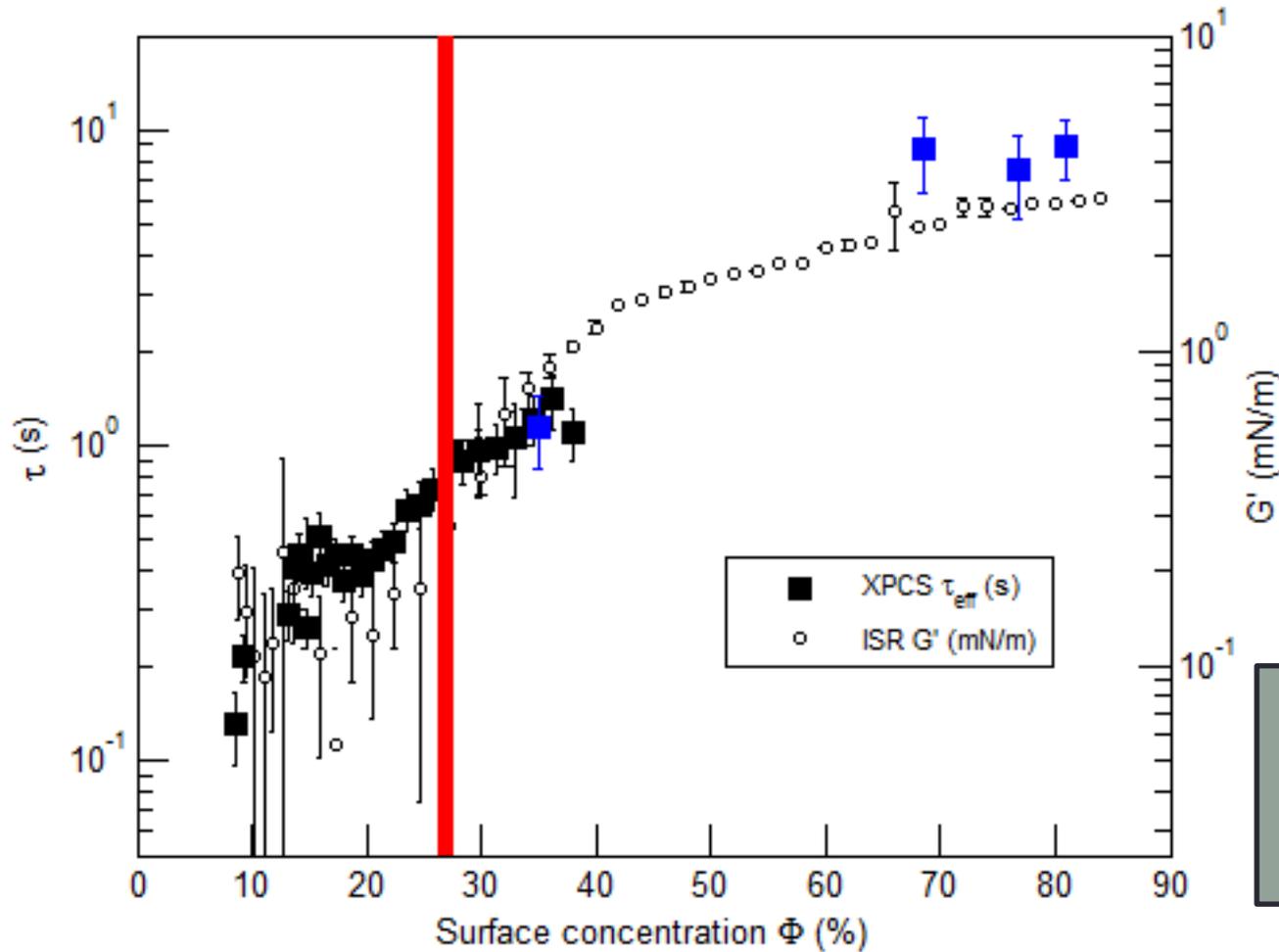
# Faster dynamics is present

The presence of dynamics faster than the exposure time is revealed by the Debye-Waller decrease of the contrast

$$\beta = \beta_0 \exp\left(-\frac{q_{\parallel}^2 \langle u^2 \rangle}{3}\right)$$



# Relaxation time $\tau$ $\longleftrightarrow$ shear modulus $G'$



This is predicted by models ascribing XPCS dynamics to stress-relaxation (e.g. Bouchaud – Pitard)

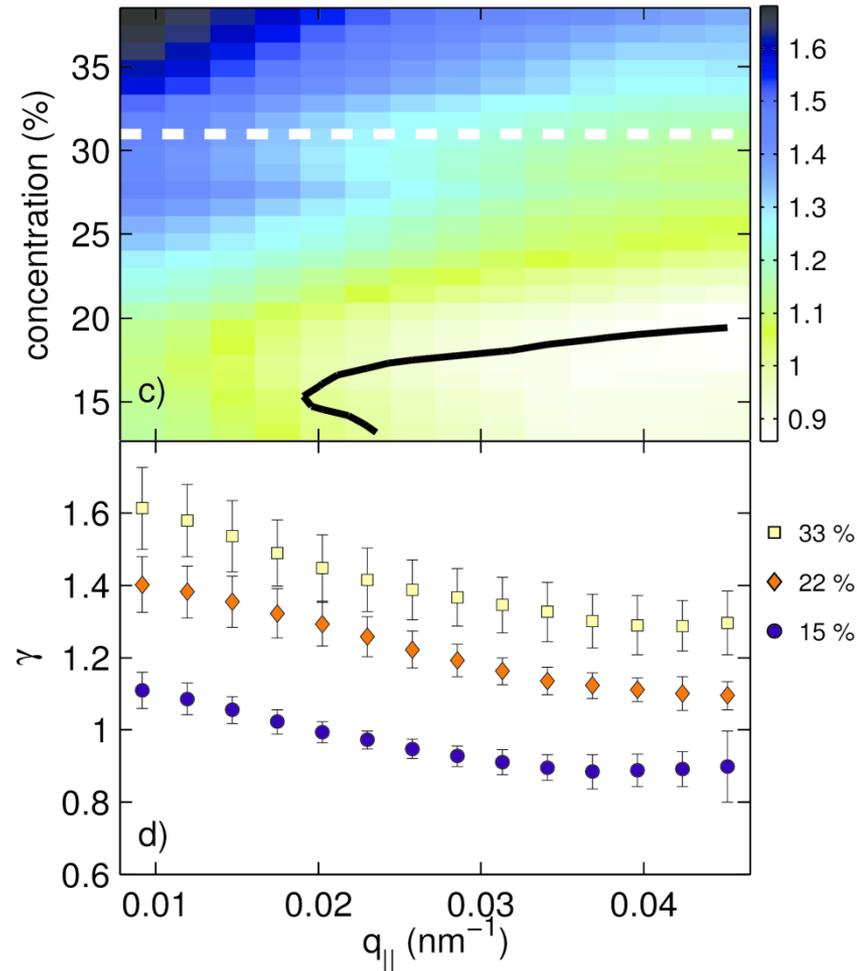
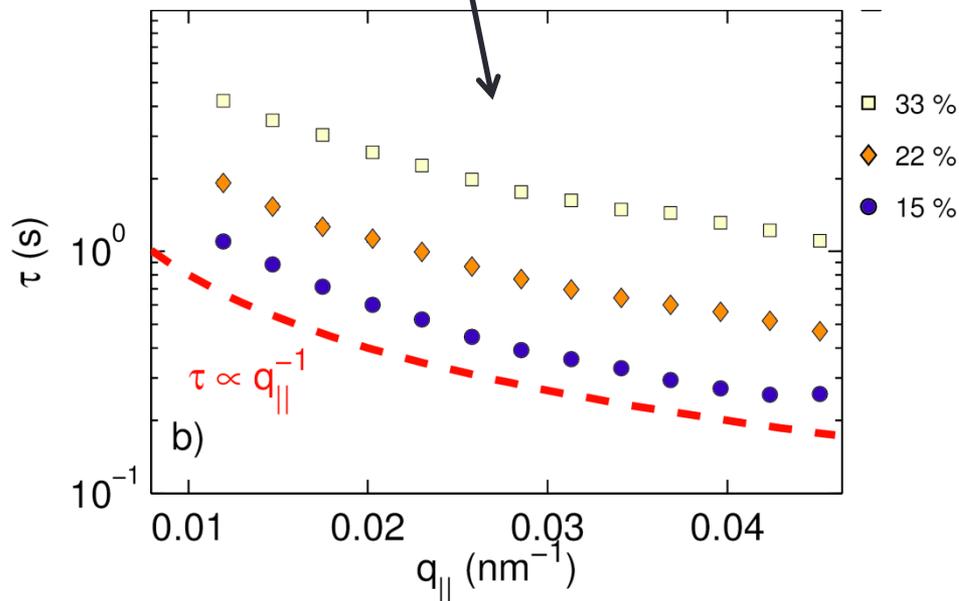
Example of Fluctuation-Mechanical Response Relation

$\tau$  and  $G'$  increase in the same way below and above the percolation threshold

# Non-brownian dynamics

Several models have been proposed to explain this behaviour  $\neq 1$

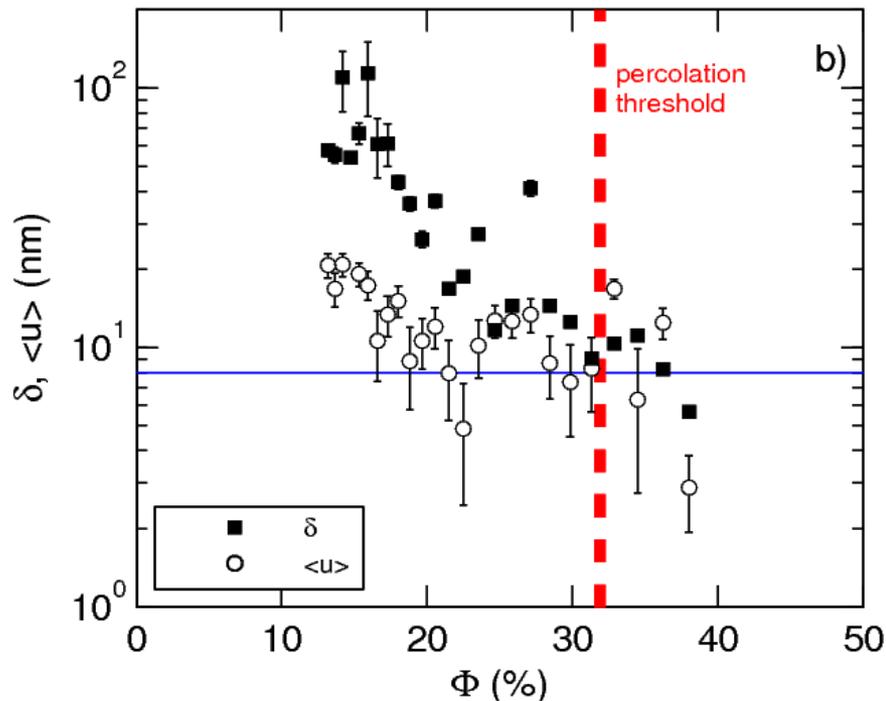
Our data seems to be well described, in particular, by a model proposed by Duri and Cipelletti:  
 $\tau \propto q_{||}^{-1}$   
 the compressed shape is due to rare, intermittent rearrangements



Also found in ageing laponite suspension - R. Angelini et al. **Soft Matter** 9, 10955 (2013)  
 metallic glasses - B. Ruta et al. **Phys. Rev. Lett.** 109, 165701 (2012)

# Intermittent rearranging events

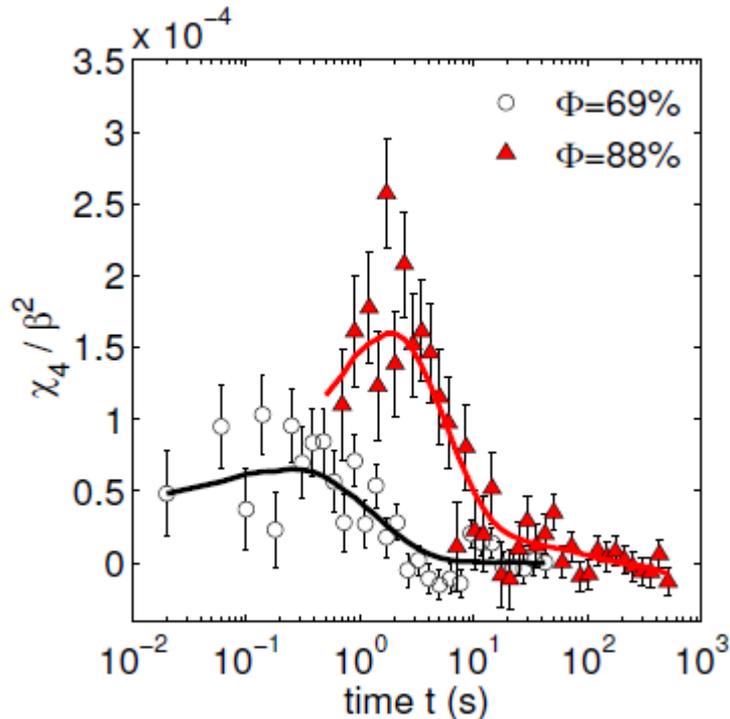
- Slow dynamics  $\rightarrow$  instantaneous, random discrete rearrangements = long range elastic deformation of the gel under the action of dipolar stresses.
- The model predicts  $\tau \propto q^{-1}$  and a  $q$ -dependent  $\gamma$ : **this holds at all concentrations**. From the model we extract the **characteristic length scale of the random rearrangements  $\delta$**  (building a master curve).



Two different regimes exist, above and below the percolation threshold: below it, the length scales  $\langle u \rangle$  and  $\delta$  of the fast and slow dynamics decouple, while above it they **become equal** within experimental accuracy, and also comparable to the size of the single nanoparticle.

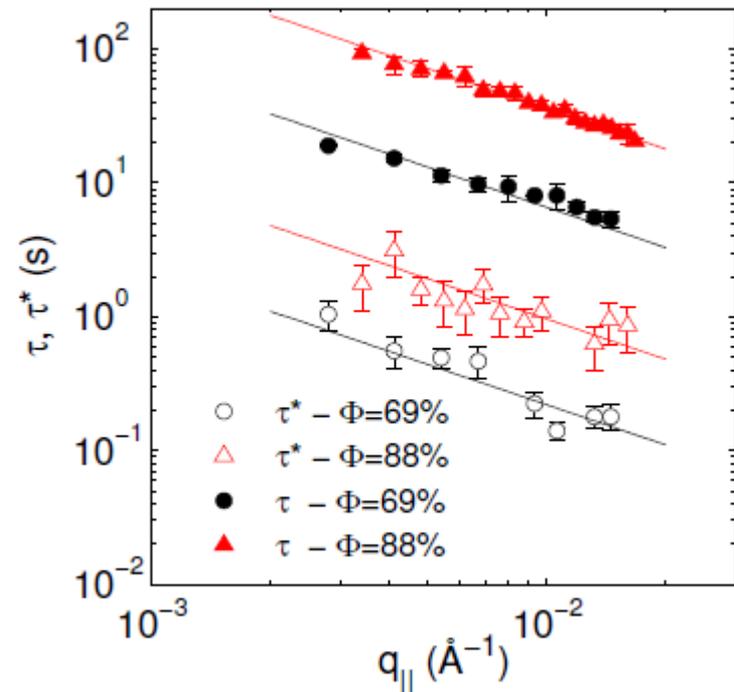
# Dynamical heterogeneities

Accessed by means of fourth order correlation functions:



At higher  $\Phi$ :

- Increase in peak - larger Cooperative Rearranging regions
- Slower dynamics



- $\tau^* \propto q^{-1}$
- Faster dynamics than  $\tau$

# Conclusions

We were able to characterize the formation of a two-dimensional gel from different, complementary perspectives:

- Structural
  - finding analogies with **2D percolation theory**
- Rheological
  - measuring the **elastic response** commonly found in gels, **even at low concentration**
  - finding interesting analogies in the concentration dependence of the shear modulus, compared to the one of the internal dynamics
- Dynamical
  - observing **non-brownian** dynamics taking place on two different spatial scales, that converge as the percolation threshold is reached
  - characterizing the **dynamical heterogeneities** of the sample

# Thank you!

- **Structure of the gel network**

- D. Orsi, A. Vezzani, R. Burioni, A. Pucci, G. Ruggeri, and L. Cristofolini, **Colloids Surfaces A Physicochem. Eng. Asp.** 441, 912 (2014)

- **Mechanical response**

- D. Orsi, G. Baldi, P. Cicutta, and L. Cristofolini, **Colloids Surfaces A Physicochem. Eng. Asp.** 413, 71 (2012).

- **Internal dynamics**

- D. Orsi, L. Cristofolini, G. Baldi, and A. Madsen, **Physical Review Letters** 108, 105701 (2012).
- D. Orsi, B. Ruta, Y. Chushkin, A. Pucci, G. Ruggeri, G. Baldi, T. Rimoldi and L. Cristofolini, **submitted to Physical Review E**

# Intermittent rearranging events

$$g_2(\tau, q) - 1 = \sum_{n=0}^{\infty} P_{\tau}(n) h(n, q)$$

- Slow dynamics  $\rightarrow$  instantaneous discrete events occurring with **Poissonian statistics**
- The degree of correlation related to  $n$  events at a given  $q$  is given by the Fourier transform of the probability distribution function (PDF) of the particle displacements, assuming **long range elastic deformation of the gel under the action of dipolar stresses**  $h(n, q) = \exp[-(qn\delta)^{1.5}]$

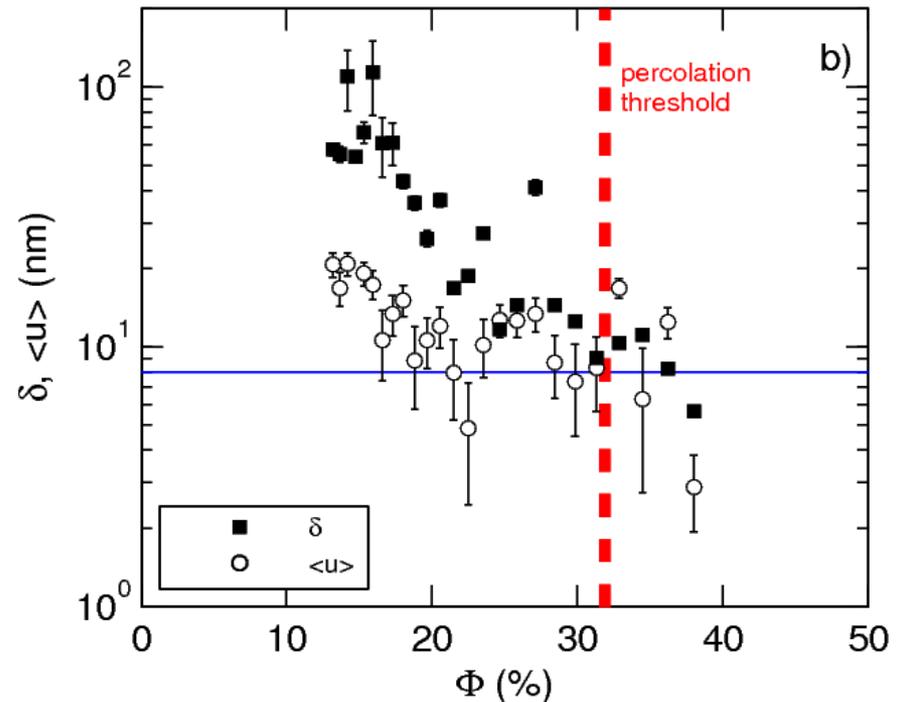
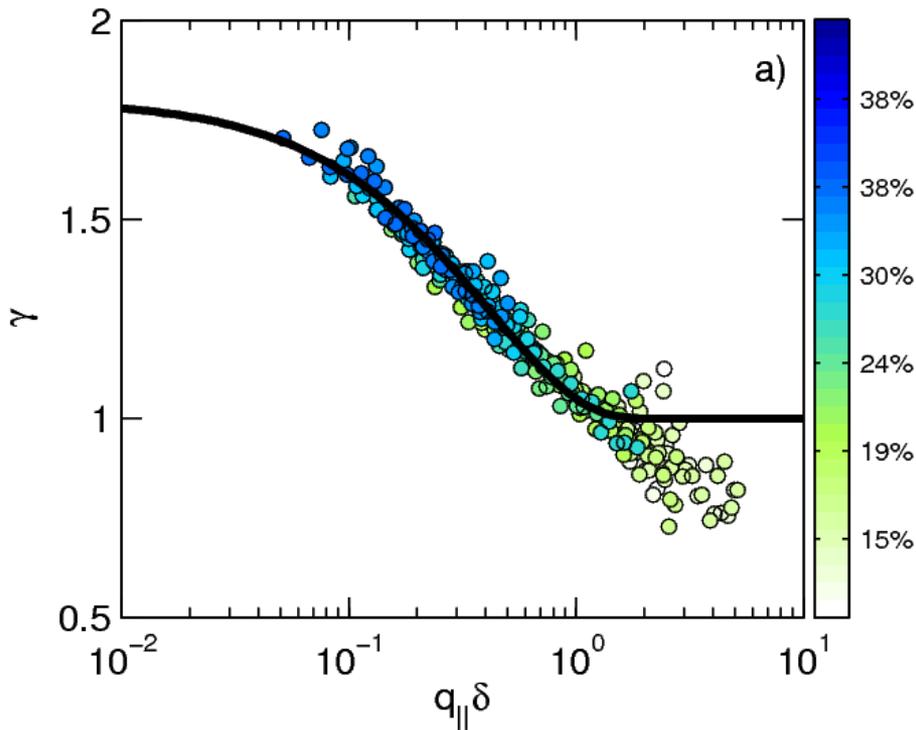
Since  $PDF \propto \Delta R^{-(1.5+1)}$  for dipolar stresses randomly scattered in space

The resulting correlation functions are fitted with stretched exponentials to retrieve the  $q$  dependence of the parameters

$$q\delta \rightarrow 0: \gamma = 1.5 \qquad q\delta \gg 1: \gamma = 1$$

# Intermittent rearranging events

The dynamics measured above and below the percolation threshold can be described within the same framework! Therefore, we build a mastercurve.



Two different regimes exist, above and below the percolation threshold: below it, the length scales  $\langle u \rangle$  and  $\delta$  of the fast and slow dynamics decouple, while above it they **become equal** within experimental accuracy, and also comparable to the size of the single nanoparticle.